

**High-Spin Molecules: Hexanuclear
[Mn₆O₄Cl₄(Me₂dbm)₆] (Me₂dbmH =
4,4'-Dimethyldibenzoylmethane) with a Near
Tetrahedral [Mn₆O₄Cl₄]⁶⁺ Core and a S = 12
Ground State**

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An interesting sub-area of transition-metal cluster chemistry is the small but growing family of molecules that, in their ground states, have unusually large numbers of unpaired electrons. Molecular clusters with spin (*S*) values as high as double figures have been discovered, with the highest value to date being *S* ≈ 16½ for one of the clusters in a sample of cocrystallized Fe₁₇ and Fe₁₉ species,² but examples with *S* ≥ 8 are nevertheless very rare.^{3,4} The study of such molecules has shown that the high spin value is a result of the presence of (at least some) ferromagnetic exchange interactions between the metal ions and/or spin frustration effects arising from the presence in certain M_x topologies of competing exchange interactions which prevent (frustrate) the preferred spin alignments that would otherwise normally yield low-spin species.

The study of high-spin molecules has taken on additional importance in recent years as it has been realized that a fairly large *S* value is a necessary (but not sufficient) property for molecules to exhibit the new magnetic phenomenon of single-molecule magnetism, i.e., the ability of material composed of discrete, (magnetically) noninteracting molecules to be magnetized by an external magnetic field below a critical blocking temperature (*T*_B).^{3a,5,6} The importance of a high *S* value in such nanoscale magnets has stimulated a search for new examples of species with this property. We herein report the preparation of a new hexanuclear Mn^{III} cluster with an aesthetically pleasing [Mn₆O₄-

Cl₄]⁶⁺ core and show that it possesses a *S* = 12 ground state, one of the highest yet identified.

Treatment of [Mn₄O₂(O₂CMe)₆(py)₂(Me₂dbm)₂] (**1**) (Me₂dbmH = 4,4'-dimethyldibenzoylmethane), prepared as for the dbm analogue,⁷ in CH₂Cl₂ with 6 equiv of Me₃SiCl and 2 equiv of Me₂dbmH gave a brown solution from which was isolated crude [MnCl(Me₂dbm)₂] (**2**) on addition of Et₂O. Redissolution in MeCN:CH₂Cl₂ (1:1) followed by standing and slight concentration at room temperature over ~3 weeks gave black crystals of [Mn₆O₄Cl₄(Me₂dbm)₆]·3CH₂Cl₂ (3·3CH₂Cl₂), together with some white powder. The latter was removed by filtration and washing of the black crystals with EtOH. Complex **3** can also be obtained by dissolving purified **2** in MeCN:CH₂Cl₂ (1:1), and subsequent treatment as above, suggesting the formation of **3** to involve a slow hydrolysis of the mononuclear species.⁸ Nonoptimized yields up to 14% have been obtained to date.⁸

The structure⁹ of **3** (Figure 1) consists of a (Mn^{III})₆ octahedron with four nonadjacent faces bridged by the μ₃-O²⁻ ions and the other four faces by μ₃-Cl⁻ ions. Six-coordinate, approximately octahedral geometry at each metal is completed by a chelating Me₂dbm group. The cluster has virtual *T_d* symmetry. As expected for high-spin, octahedral Mn^{III}, there is a Jahn–Teller (JT) distortion, taking the form of an axial elongation of the two trans Mn–Cl bonds, making them unusually long (2.618(3)–2.692(3) Å). In contrast, the Mn–O²⁻ (1.876(4)–1.899(5) Å) and Mn–O(Me₂dbm) (1.903(5)–1.925(5) Å) bond lengths are as expected. As a result of (i) the long Mn–Cl⁻ versus short Mn–O²⁻ bonds and (ii) the near trigonal planar geometry at the latter (sum-of-angles (soa) ≈ 349°) compared with marked trigonal pyramidal geometry at the former (soa ≈ 223°), the [Mn₆O₄Cl₄]⁶⁺ core is a near tetrahedron with a Cl⁻ at each vertex, a Mn at the midpoint of each edge, and a O²⁻ bridging each face. Although many [M₆(μ₃-X)₈] face-capped metal octahedra are known,¹⁰ only a relative few contain two types of X group, e.g., the [Ti₆O₆Cl₂],¹¹ [Ti₆Te₆O₂],¹² and [Re₆Y_xZ_{8-x}] (*x* = 5, Y = S or Se, Z = Cl; *x* = 6, Y = S, Z = Cl)¹³ cores. Only the [Ti₆(μ₃-O)₄(μ₃-Cl)₄] core of [(C₅H₅Me)₆Ti₆O₄Cl₄]¹¹ contains, like **3**, four O²⁻ and 4 Cl⁻ bridges, but the structure does not approximate to a tetrahedron. [Mn^{III}₆X₈] species have been unknown to date, although a [Mn₆(μ₃-O)₄(μ₃-Cl)₄]⁴⁺ unit as found in **3** but at the 2Mn^{II}, 4Mn^{III} level is also a recognizable fragment within the higher nuclearity cluster [Mn₁₀O₄Cl₁₂(biphen)₄]⁴⁻ (biphen = 2,2'-biphenoxide).¹⁴

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(8) The reaction is more complicated than this statement suggests, with [Mn(Me₂dbm)₃] identified in the filtrate; the isolation of pure **3** is undoubtedly due to its low solubility in this solvent mixture. Isolated solid is soluble in CH₂Cl₂ and CHCl₃. ¹H NMR spectra in these solvents suggest the structure is retained on dissolution.

(9) Dried solid analyzed as 3·0.4CH₂Cl₂. Anal. Calcd (found) for C_{102.4}H_{90.8}O₁₆Cl_{4.8}Mn₆: C, 59.21 (59.28); H, 4.49 (4.41). Crystal data for 3·3CH₂Cl₂: monoclinic; *P*2₁/*c*; *a* = 17.172(2) Å, *b* = 18.302(2) Å, *c* = 34.534(4) Å; β = 100.36(1)°; *Z* = 4; *V* = 10677 Å³; *d*_{calcd} = 1.430 g cm⁻³; *T* = -171 °C. The structure was solved using MULTAN and refined on *F* to *R*(*R*_w) = 5.80 (5.73) using 9989 unique reflections with *F* > 2.33σ(*F*). All non-hydrogen atoms were refined anisotropically, and all hydrogen atoms were included as fixed-atom contributors at calculated positions, except for the disordered CH₂-Cl₂ molecules. Electronic spectrum in CH₂Cl₂, λ_{max}/nm (ε_m/L mol⁻¹ cm⁻¹): 456 (6380), 486 (5000), 540 (2720), 584 1930.

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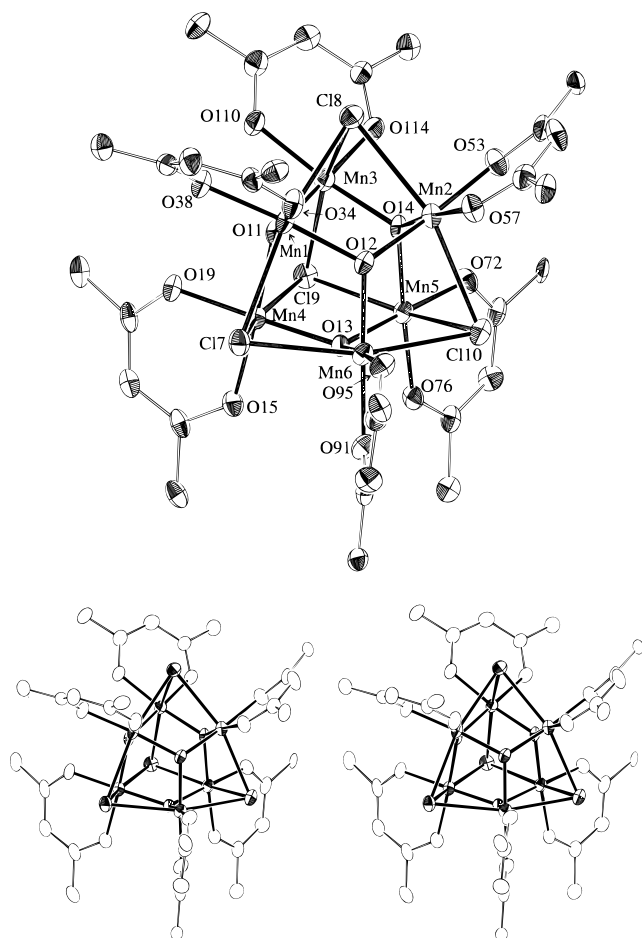


Figure 1. Labeled ORTEP plot and stereopair of $[\text{Mn}_6\text{O}_4\text{Cl}_4(\text{Me}_2\text{dbm})_6]$ (**3**). For clarity, only the *ipso*-carbon atom of each C_6H_4 -*p*-Me ring is shown. Selected distances (Å): Mn1–Mn2, 3.195(2); Mn1–Mn3, 3.207(2); Mn1–Mn4, 3.203(2); Mn1–Mn6, 3.193(2); Mn2–Mn3, 3.199(2); Mn2–Mn5, 3.210(2); Mn2–Mn6, 3.219(2); Mn3–Mn4, 3.204(2); Mn3–Mn5, 3.189(2); Mn4–Mn5, 3.205(2); Mn4–Mn6, 3.199(2); Mn5–Mn6, 3.221(2); Mn1–Mn5, 4.526(2); Mn2–Mn4, 4.534(2); Mn3–Mn6, 4.532(2).

Variable-temperature, magnetic susceptibility data were collected on powdered **3**·0.4 CH_2Cl_2 in the range of 2.00–320 K. The $\mu_{\text{eff}}/\mu_{\text{B}}$ ($\chi_{\text{m}}T/\text{cm}^3 \text{K mol}^{-1}$) values slowly increase from 16.01 (32.04) at 320 K to a maximum of 24.27 (73.63) at 30.0 K and then decrease to 13.69 (23.43) at 2.00 K. The maxima may be compared with 24.99 (78.0) for a $S = 12$ system with $g = 2.00$. This ground-state value was confirmed by fitting of reduced magnetization ($M/N\mu_{\text{B}}$) vs H/T data collected in the 0.500–50.0 kG range and temperatures down to 2.00 K (Figure 2). The $M/N\mu_{\text{B}}$ value saturates at 23.02, near to the value of 24 expected for $S = 12$ if $g = 2.0$. The various isofield lines are virtually superimposed, indicating essentially no zero field splitting (ZFS) within the ground state. The data were fit (solid line in Figure 2) to a Brillouin function for an isolated $S = 12$ state with $g = 1.936$ and $D = 0.0 \text{ cm}^{-1}$, confirming a high-spin ground state with little or no magnetic anisotropy, as anticipated from the

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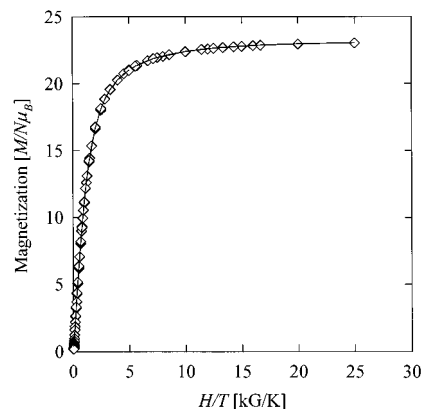


Figure 2. Plot of reduced magnetization ($M/N\mu_{\text{B}}$) vs H/T for $[\text{Mn}_6\text{O}_4\text{Cl}_4(\text{Me}_2\text{dbm})_6] \cdot 0.4\text{CH}_2\text{Cl}_2$, collected in the 2.00–15.0 K and 0.50–50 kG ranges. The solid line is a fit of the data to the Brillouin function for a $S = 12$ complex with $g = 1.936$.

virtual T_d symmetry. To obtain the pairwise exchange parameters, the μ_{eff} versus T data were fit to the expression derived for a $(\text{Mn}^{\text{III}})_6$ octahedron using the Kambe vector coupling method,¹⁵ the van Vleck equation, and exchange parameters J_{cis} and J_{trans} . The spin Hamiltonian is given in eq 1

$$\hat{H} = -2J_{\text{cis}}(\hat{S}_1\hat{S}_2 + \hat{S}_1\hat{S}_3 + \hat{S}_1\hat{S}_4 + \hat{S}_1\hat{S}_6 + \hat{S}_2\hat{S}_3 + \hat{S}_2\hat{S}_5 + \hat{S}_2\hat{S}_6 + \hat{S}_3\hat{S}_4 + \hat{S}_3\hat{S}_5 + \hat{S}_4\hat{S}_5 + \hat{S}_4\hat{S}_6 + \hat{S}_5\hat{S}_6) - 2J_{\text{trans}}(\hat{S}_1\hat{S}_5 + \hat{S}_2\hat{S}_4 + \hat{S}_3\hat{S}_6) \quad (1)$$

(where S_i is the spin of metal Mn_i), which can be transformed into the equivalent form in eq 2 by the substitutions $\hat{S}_A = \hat{S}_1 + \hat{S}_5$, $\hat{S}_B = \hat{S}_2 + \hat{S}_4$, $\hat{S}_C = \hat{S}_3 + \hat{S}_6$, and $\hat{S}_T = \hat{S}_A + \hat{S}_B + \hat{S}_C$, where \hat{S}_T is the resultant spin of the complete molecule. The energies,

$$\hat{H} = -J_{\text{cis}}(\hat{S}_T^2 - \hat{S}_A^2 - \hat{S}_B^2 - \hat{S}_C^2) - J_{\text{trans}}(\hat{S}_A^2 + \hat{S}_B^2 + \hat{S}_C^2 - \hat{S}_1^2 - \hat{S}_2^2 - \hat{S}_3^2 - \hat{S}_4^2 - \hat{S}_5^2 - \hat{S}_6^2) \quad (2)$$

$E(S_T)$, of each S_T are given by eq 3, where constant terms have been omitted. An excellent fit was obtained with $J_{\text{cis}} = +8.6$

$$E(S_T) = -J_{\text{cis}}[S_T(S_T + 1) - S_A(S_A + 1) - S_B(S_B + 1) - S_C(S_C + 1)] - J_{\text{trans}}[S_A(S_A + 1) + S_B(S_B + 1) + S_C(S_C + 1)] \quad (3)$$

cm^{-1} , $J_{\text{trans}} = 0 \text{ cm}^{-1}$, and $g = 1.965$ with TIP held constant at $1200 \times 10^{-6} \text{ cm}^3 \text{K mol}^{-1}$. These J values indicate a well-isolated $S = 12$ ground state separated by 138 cm^{-1} from the first excited state.

The above results establish **3** as a new member of the family of high-spin molecules and represent the first time that a discrete $(\text{Mn}^{\text{III}})_6$ octahedron has been prepared. We believe it to be a prototype of a large new family of related complexes, and attempts to prepare the Br^- and other analogues are currently in progress.

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Supporting Information Available: Crystallographic data collection and refinement details and listings of atomic coordinates and thermal parameters for **3**·3 CH_2Cl_2 (21 pages, print/PDF). See any current masthead page for ordering information and Web access instructions.

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